

NumRu::GPhys::EP_Flux 数理ドキュメント

地球流体電脳倶楽部

平成17年2月19日

目次

第1章	はじめに	2
第2章	NumRu::GPhys::EP_Flux で計算される緒量	3
2.1	系の設定	3
2.2	EP フラックス	4
2.3	残差循環	5
2.4	平均東西流の式	5
2.5	子午面上の発散演算子	5
2.6	質量流線関数	5
2.7	変数変換	6
付録A	プリミティブ方程式系と変形オイラー平均の復習	7
1.1	球面上の対数圧力座標系におけるプリミティブ方程式	7
1.2	オイラー平均方程式系	8
1.3	変形オイラー平均方程式系	13

第1章 はじめに

NumRu::GPhys::EP_Flux は Eliassen-Palm フラックス (EP フラックス) および残差循環を計算するメソッドを集めたモジュールである。現状では、鉛直座標として対数圧力座標を用いた球座標系におけるプリミティブ方程式 (準地衡風近似をしない) EP フラックスのためのメソッドだけが用意されている。将来的には Plumb フラックスや Takaya-Nakamura フラックスを計算するメソッドもサポートする予定である。本ドキュメントでは NumRu::GPhys::EP_Flux で使用される数式の解説と各メソッドの概説を行う。なお、NumRu::GPhys::EP_Flux では、微分演算のために、別モジュール NumRu::Derivative および NumRu::GPhys::Derivative で定義されるメソッドを使用している。微分演算メソッドに関する詳細はそれぞれのモジュールのドキュメントを参照されたい。

第2章 NumRu::GPhys::EP_Flux で計算される緒量

本章では NumRu::GPhys::EP_Flux で定義される緒量の解説を行う。数理モデルは Andrews *et al.*(1987) の第3章に基づく。

2.1 系の設定

球面上の大気を考える。大気の厚さは水平方向の広がりには比べ薄く、鉛直方向に静水圧平衡が成り立つものとする。緯度経度座標系を用い、経度 λ 軸を東向き、緯度 ϕ 軸を北向きに正をとる。鉛直座標には対数圧力座標 z^*

$$z^* = -H \ln(p/p_s), \quad H = \frac{R_d T_s}{g_0} \quad (2.1)$$

を用いる。ここで H はスケールハイト、 R_d は乾燥空気の気体定数 (普遍気体定数を R , 乾燥空気の分子量を w とすると $R_d = R/w$), T_s は標準参照温度 (定数), g_0 は地表面における重力加速度 (定数), p は圧力, p_s は参照圧力である。 p_s として地表面圧力の代表値 (定数) を用いる。

2.2 EP フラックス

本モジュールでは惑星半径と後述の ρ_s で規格化した EP フラックス (以降, 規格化した EP フラックス) を計算, 出力する. 規格化した EP フラックスは

$$\hat{F}_\phi \equiv \sigma \cos \phi \left(\frac{\partial \bar{u}}{\partial z^*} \frac{\overline{v'\theta'}}{\frac{\partial \bar{\theta}}{\partial z^*}} - \overline{u'v'} \right), \quad (2.2a)$$

$$\hat{F}_{z^*} \equiv \sigma \cos \phi \left(\left[f - \frac{1}{a \cos \phi} \frac{\partial \bar{u} \cos \phi}{\partial \phi} \right] \frac{\overline{v'\theta'}}{\frac{\partial \bar{\theta}}{\partial z^*}} - \overline{u'w'} \right) \quad (2.2b)$$

と定義される. ここで \hat{F}_ϕ , \hat{F}_{z^*} はそれぞれ規格化された EP フラックスの ϕ 成分, z^* 成分である. $\bar{\bullet}$ は東西オイラー平均量, \bullet' は東西オイラー平均量からのずれを表す. u, v, w はそれぞれ東西風速, 南北風速, 対数圧力速度で

$$(u, v, w) \equiv \left(a \cos \phi \frac{d\lambda}{dt}, a \frac{d\phi}{dt}, \frac{dz^*}{dt} \right)$$

と定義される. θ は温位, a は惑星半径 (定数) である. σ は

$$\sigma \equiv \frac{\rho_0}{\rho_s} = \exp \left(\frac{-z^*}{H} \right), \quad (2.3)$$

である. ただし, ρ_0 は基本場の密度で

$$\rho_0(z^*) \equiv \rho_s e^{-z^*/H}, \quad \rho_s \equiv p_s / RT_s$$

である. f はコリオリパラメータで

$$f = 2\Omega \sin \phi = \frac{4\pi}{T_{rot}} \sin \phi \quad (2.4)$$

と定義される. Ω は自転角速度, T_{rot} は自転周期である. 本モジュールでは, 自転角速度を変更するためには T_{rot} の値を与える仕様になっている.

一方, Andrews *et al.* (1987) で示されている EP フラックスは次のように定義される

$$F_\phi = \rho_0 a \cos \phi \left(\frac{\partial \bar{u}}{\partial z^*} \frac{\overline{v'\theta'}}{\frac{\partial \bar{\theta}}{\partial z^*}} - \overline{u'v'} \right) \quad (2.5a)$$

$$F_{z^*} = \rho_0 a \cos \phi \left(\left[f - \frac{\partial \bar{u} \cos \phi}{\partial \phi} \right] \frac{\overline{v'\theta'}}{\frac{\partial \bar{\theta}}{\partial z^*}} - \overline{u'w'} \right). \quad (2.5b)$$

ここで F_ϕ , F_{z^*} はそれぞれ EP フラックスの ϕ 成分, z^* 成分である. F_y, F_z と \hat{F}_y, \hat{F}_z は以下のように関係付けられる.

$$(F_y, F_z) = a \rho_s (\hat{F}_y, \hat{F}_z) \quad (2.6)$$

2.3 残差循環

残差循環 $(0, \bar{v}^*, \bar{w}^*)$ は以下の形で定義される.

$$\begin{aligned}\bar{v}^* &\equiv \bar{v} - \frac{1}{\rho_0} \frac{\partial}{\partial z^*} \left(\rho_0 \frac{\overline{v'\theta'}}{\frac{\partial \bar{\theta}}{\partial z^*}} \right) \\ &= \bar{v} - \frac{1}{\sigma} \frac{\partial}{\partial z^*} \left(\sigma \frac{\overline{v'\theta'}}{\frac{\partial \bar{\theta}}{\partial z^*}} \right)\end{aligned}\quad (2.7a)$$

$$\bar{w}^* \equiv \bar{w} + \frac{1}{a \cos \phi} \frac{\partial}{\partial \phi} \left(\cos \phi \frac{\overline{v'\theta'}}{\frac{\partial \bar{\theta}}{\partial z^*}} \right)\quad (2.7b)$$

2.4 平均東西流の式

規格化した EP フラックスを用いると, TEM 系における u の式は以下のようになる.

$$\frac{\partial \bar{u}}{\partial t} + \bar{v}^* \left[\frac{1}{a \cos \phi} \frac{\partial}{\partial \phi} (\bar{u} \cos \phi) - f \right] + \bar{w}^* \frac{\partial \bar{u}}{\partial z^*} - \bar{X} = \frac{1}{\sigma \cos \phi} \nabla \cdot \hat{\mathbf{F}}.\quad (2.8)$$

2.5 子午面上の発散演算子

子午面における発散演算子は, \mathbf{F} を任意のベクトルした時に以下の形で定義される.

$$\nabla \cdot \mathbf{F} = \frac{1}{a \cos \phi} \frac{\partial (\cos \phi F_\phi)}{\partial \phi} + \frac{\partial F_{z^*}}{\partial z^*}\quad (2.9)$$

2.6 質量流線関数

残差循環の質量流線関数 Ψ^* を

$$\sigma \bar{v}^* = -g \frac{1}{2\pi a \cos \phi} \frac{\partial \Psi^*}{\partial z^*},\quad (2.10a)$$

$$\sigma \bar{w}^* = g \frac{1}{2\pi a^2 \cos \phi} \frac{\partial \Psi^*}{\partial \phi}\quad (2.10b)$$

と定義する. 上式を積分して Ψ^* を求めるために, 本モジュールでは (2.1) を使用して対数圧力座標 (z^*) 系から圧力座標 (p) 系へ

$$\frac{\partial}{\partial z^*} \Psi^* = -\frac{p}{H} \frac{\partial}{\partial p} \Psi^* \quad (2.11)$$

と変換し, 大気上端 ($p = 0$) において $\Psi^* = 0$ として積分し

$$\Psi^*(\theta, p) = \frac{2\pi a \cos \phi}{g} \int_0^p \bar{v}^* dp \quad (2.12)$$

と質量流線関数を導いている.

2.7 変数変換

EP_Flux モジュールでは与えられたデータに応じて変数変換を施す場合がある. その変換は以下のように行う.

入力されるデータの鉛直軸が気圧軸であった場合, 以下の関係式を用いて高度軸に変換し, 計算を行う.

$$z^* = -H \log \left(\frac{p}{p_{00}} \right), \quad (2.13a)$$

$$p = p_{00} \exp \left(-\frac{z^*}{H} \right) \quad (2.13b)$$

ここで p は圧力, p_{00} は地表面参考気圧 (定数) である.

入力が θ や w でなく, 気温 T , 圧力「速度」 $\omega \equiv Dp/Dt$ の場合はそれぞれを元に w, θ を求める必要がある. 本モジュールでは以下の式を用いて w, θ を求める.

$$w = -\omega H/p \quad (2.14)$$

$$\theta = T \left(\frac{p_{00}}{p} \right)^\kappa, \kappa = R/C_p \quad (2.15)$$

ここで R, C_p はそれぞれ乾燥空気の気体定数および定圧比熱である.

付録A プリミティブ方程式系と変形オイラー平均の復習

本章では変形オイラー平均方程式系と EP フラックス および残差循環の関係を確認する。まず対数圧力座標系を用いた球面上の 3 次元プリミティブ方程式を提示する。次いでそのオイラー平均および変形オイラー平均方程式系を導出する。最後に変形オイラー平均方程式に基づき EP フラックスおよび残差循環を定義する。

1.1 球面上の対数圧力座標系におけるプリミティブ方程式

球面上の対数圧力座標系におけるプリミティブ方程式は以下の通りである。ここでは Andrews *et al.* (1987) の (3.1.3) 式を参考にした。

$$\frac{du}{dt} - \left(f + \frac{u \tan \phi}{a} \right) v + \frac{1}{a \cos \phi} \frac{\partial \Phi}{\partial \lambda} = X, \quad (\text{A.1a})$$

$$\frac{dv}{dt} + \left(f + \frac{u \tan \phi}{a} \right) u + \frac{1}{a} \frac{\partial \Phi}{\partial \phi} = Y, \quad (\text{A.1b})$$

$$\frac{\partial \Phi}{\partial z^*} = \frac{R\theta e^{-\kappa z^*/H}}{H}, \quad (\text{A.1c})$$

$$\frac{1}{a \cos \phi} \left[\frac{\partial u}{\partial \lambda} + \left(\frac{\partial v \cos \phi}{\partial \phi} \right) \right] + \frac{1}{\rho_0} \frac{\partial}{\partial z^*} (\rho_0 w) = 0, \quad (\text{A.1d})$$

$$\frac{d\theta}{dt} = Q, \quad (\text{A.1e})$$

ここで Φ はジオポテンシャルハイト, X, Y はそれぞれ外力の λ 成分 と ϕ 成分, $\kappa = R_d/c_p$ (c_p は等圧比熱) である. Q は非断熱加熱項で,

$$Q = \frac{J}{C_p} e^{\kappa z^*/H}$$

である. J は単位質量あたりの非断熱加熱率である. ここで明記した以外の変数の定義については第 2.1 節, 第 2.2 節 を参照のこと.

1.2 オイラー平均方程式系

ある物理量 A について, ϕ, z^*, t を固定して東西方向にとった平均

$$\bar{A}(\phi, z^*, t) \equiv \frac{1}{2\pi} \int_0^{2\pi} A(\lambda, \phi, z^*, t) d\lambda \quad (\text{A.2})$$

をオイラー平均と呼ぶ. オイラー平均からのずれを A' とすると

$$A' = A - \bar{A} \quad (\text{A.3})$$

である. 定義により, $\bar{A}' = 0, \partial \bar{A}' / \partial \lambda = 0$ となる.

(A.1) 中の各量をオイラー平均とそこからのずれに分けて書くと

$$\begin{aligned} \frac{\partial}{\partial t}(\bar{u} + u') + \frac{\bar{u} + u'}{a \cos \phi} \frac{\partial}{\partial \lambda}(\bar{u} + u') + \frac{\bar{v} + v'}{a} \frac{\partial}{\partial \phi}(\bar{u} + u') + (\bar{w} + w') \frac{\partial}{\partial z^*}(\bar{u} + u') \\ - \left[f + \frac{\tan \phi}{a}(\bar{u} + u') \right] (\bar{v} + v') + \frac{1}{a \cos \phi} \frac{\partial}{\partial \lambda}(\bar{\Phi} + \Phi') = \bar{X} + X', \end{aligned} \quad (\text{A.4a})$$

$$\begin{aligned} \frac{\partial}{\partial t}(\bar{v} + v') + \frac{\bar{u} + u'}{a \cos \phi} \frac{\partial}{\partial \lambda}(\bar{v} + v') + \frac{\bar{v} + v'}{a} \frac{\partial}{\partial \phi}(\bar{v} + v') + (\bar{w} + w') \frac{\partial}{\partial z^*}(\bar{v} + v') \\ + \left[f + \frac{\tan \phi}{a}(\bar{u} + u') \right] (\bar{u} + u') + \frac{1}{a} \frac{\partial}{\partial \phi}(\bar{\Phi} + \Phi') = \bar{Y} + Y', \end{aligned} \quad (\text{A.4b})$$

$$\frac{\partial}{\partial z^*}(\bar{\Phi} + \Phi') = \frac{Re^{-\kappa z^*/H}}{H}(\bar{\theta} + \theta'), \quad (\text{A.4c})$$

$$\frac{1}{a \cos \phi} \left[\frac{\partial}{\partial \lambda}(\bar{u} + u') + \frac{\partial}{\partial \phi} \{(\bar{v} + v') \cos \phi\} \right] + \frac{1}{\rho_0} \frac{\partial}{\partial z^*}[\rho_0(\bar{w} + w')] = 0, \quad (\text{A.4d})$$

$$\begin{aligned} \frac{\partial}{\partial t}(\bar{\theta} + \theta') + \frac{\bar{u} + u'}{a \cos \phi} \frac{\partial}{\partial \lambda}(\bar{\theta} + \theta') + \frac{\bar{v} + v'}{a} \frac{\partial}{\partial \phi}(\bar{\theta} + \theta') + (\bar{w} + w') \frac{\partial}{\partial z^*}(\bar{\theta} + \theta') \\ = \bar{Q} + Q' \end{aligned} \quad (\text{A.4e})$$

となる. 上記を変形して, 左辺に平均量と平均量同士の積の項を, 右辺にそれ以外の項をまとめると

$$\begin{aligned}
& \frac{\partial \bar{u}}{\partial t} + \frac{\bar{u}}{a \cos \phi} \frac{\partial \bar{u}}{\partial \lambda} + \frac{\bar{v}}{a} \frac{\partial \bar{u}}{\partial \phi} + \bar{w} \frac{\partial \bar{u}}{\partial z^*} - f\bar{v} - \frac{\tan \phi}{a} \bar{u} \bar{v} + \frac{1}{a \cos \phi} \frac{\partial \bar{\Phi}}{\partial \lambda} - \bar{X} \\
&= -\frac{\partial u'}{\partial t} - \frac{\bar{u}}{a \cos \phi} \frac{\partial u'}{\partial \lambda} - \frac{u'}{a \cos \phi} \frac{\partial \bar{u}}{\partial \lambda} - \frac{u'}{a \cos \phi} \frac{\partial u'}{\partial \lambda} \\
&\quad - \frac{\bar{v}}{a} \frac{\partial u'}{\partial \phi} - \frac{v'}{a} \frac{\partial \bar{u}}{\partial \phi} - \frac{v'}{a} \frac{\partial u'}{\partial \phi} - \bar{w} \frac{\partial u'}{\partial z^*} - w' \frac{\partial \bar{u}}{\partial z^*} - w' \frac{\partial u'}{\partial z^*} + f v' \\
&\quad + \frac{\tan \phi}{a} \bar{u} v' + \frac{\tan \phi}{a} u' \bar{v} + \frac{\tan \phi}{a} u' v' - \frac{1}{a \cos \phi} \frac{\partial \Phi'}{\partial \lambda} + X', \tag{A.5a}
\end{aligned}$$

$$\begin{aligned}
& \frac{\partial \bar{v}}{\partial t} + \frac{\bar{u}}{a \cos \phi} \frac{\partial \bar{v}}{\partial \lambda} + \frac{\bar{v}}{a} \frac{\partial \bar{v}}{\partial \phi} + \bar{w} \frac{\partial \bar{v}}{\partial z^*} + f\bar{u} + \frac{\tan \phi}{a} (\bar{u})^2 + \frac{1}{a} \frac{\partial \bar{\Phi}}{\partial \phi} - \bar{Y} \\
&= -\frac{\partial v'}{\partial t} - \frac{\bar{u}}{a \cos \phi} \frac{\partial v'}{\partial \lambda} - \frac{u'}{a \cos \phi} \frac{\partial \bar{v}}{\partial \lambda} - \frac{u'}{a \cos \phi} \frac{\partial v'}{\partial \lambda} \\
&\quad - \frac{\bar{v}}{a} \frac{\partial v'}{\partial \phi} - \frac{v'}{a} \frac{\partial \bar{v}}{\partial \phi} - \frac{v'}{a} \frac{\partial v'}{\partial \phi} - \bar{w} \frac{\partial v'}{\partial z^*} - w' \frac{\partial \bar{v}}{\partial z^*} - w' \frac{\partial v'}{\partial z^*} - f u' \\
&\quad - 2 \frac{\tan \phi}{a} \bar{u} u' - \frac{\tan \phi}{a} (u')^2 - \frac{1}{a \cos \phi} \frac{\partial \Phi'}{\partial \phi} + Y', \tag{A.5b}
\end{aligned}$$

$$\frac{\partial \bar{\Phi}}{\partial z^*} - \frac{Re^{-\kappa z^*/H}}{H} \bar{\theta} = -\frac{\partial \Phi'}{\partial z^*} + \frac{Re^{-\kappa z^*/H}}{H} \theta', \tag{A.5c}$$

$$\begin{aligned}
& \frac{1}{a \cos \phi} \left[\frac{\partial \bar{u}}{\partial \lambda} + \frac{\partial}{\partial \phi} (\bar{v} \cos \phi) \right] + \frac{1}{\rho_0} \frac{\partial}{\partial z^*} (\rho_0 \bar{w}) \\
&= -\frac{1}{a \cos \phi} \left[\frac{\partial u'}{\partial \lambda} + \frac{\partial}{\partial \phi} (v' \cos \phi) \right] - \frac{1}{\rho_0} \frac{\partial}{\partial z^*} (\rho_0 w'), \tag{A.5d}
\end{aligned}$$

$$\begin{aligned}
& \frac{\partial \bar{\theta}}{\partial t} + \frac{\bar{u}}{a \cos \phi} \frac{\partial \bar{\theta}}{\partial \lambda} + \frac{\bar{v}}{a} \frac{\partial \bar{\theta}}{\partial \phi} + \bar{w} \frac{\partial \bar{\theta}}{\partial z^*} - \bar{Q} \\
&= -\frac{\partial \theta'}{\partial t} - \frac{\bar{u}}{a \cos \phi} \frac{\partial \theta'}{\partial \lambda} - \frac{u'}{a \cos \phi} \frac{\partial \bar{\theta}}{\partial \lambda} - \frac{u'}{a \cos \phi} \frac{\partial \theta'}{\partial \lambda} \\
&\quad - \frac{\bar{v}}{a} \frac{\partial \theta'}{\partial \phi} - \frac{v'}{a} \frac{\partial \bar{\theta}}{\partial \phi} - \frac{v'}{a} \frac{\partial \theta'}{\partial \phi} - \bar{w} \frac{\partial \theta'}{\partial z^*} - w' \frac{\partial \bar{\theta}}{\partial z^*} - w' \frac{\partial \theta'}{\partial z^*} + Q' \tag{A.5e}
\end{aligned}$$

と書ける. (A.5) をオイラー平均すると,

$$\begin{aligned} \frac{\partial \bar{u}}{\partial t} + \frac{1}{a} \bar{v} \frac{\partial \bar{u}}{\partial \phi} + \bar{w} \frac{\partial \bar{u}}{\partial z^*} - f \bar{v} - \frac{\tan \phi}{a} \bar{u} \bar{v} - \bar{X} \\ = -\frac{1}{a \cos \phi} \overline{u' \frac{\partial u'}{\partial \lambda}} - \frac{1}{a} \overline{v' \frac{\partial u'}{\partial \phi}} - \overline{w' \frac{\partial u'}{\partial z^*}} + \frac{\tan \phi}{a} \overline{u' v'}, \end{aligned} \quad (\text{A.6a})$$

$$\begin{aligned} \frac{\partial \bar{v}}{\partial t} + \frac{\bar{v}}{a} \frac{\partial \bar{v}}{\partial \phi} + \bar{w} \frac{\partial \bar{v}}{\partial z^*} + f \bar{u} + \frac{\tan \phi}{a} (\bar{u})^2 + \frac{1}{a} \frac{\partial \bar{\Phi}}{\partial \phi} - \bar{Y} \\ = -\frac{1}{a \cos \phi} \overline{u' \frac{\partial v'}{\partial \lambda}} - \frac{1}{a} \overline{v' \frac{\partial v'}{\partial \phi}} - \overline{w' \frac{\partial v'}{\partial z^*}} - \frac{\tan \phi}{a} \overline{u'^2}, \end{aligned} \quad (\text{A.6b})$$

$$\frac{\partial \bar{\Phi}}{\partial z^*} - \frac{Re^{-\kappa z^*/H}}{H} \bar{\theta} = 0, \quad (\text{A.6c})$$

$$\frac{1}{a \cos \phi} \left[\frac{\partial}{\partial \phi} (\bar{v} \cos \phi) \right] + \frac{1}{\rho_0} \frac{\partial}{\partial z^*} (\rho_0 \bar{w}) = 0, \quad (\text{A.6d})$$

$$\frac{\partial \bar{\theta}}{\partial t} + \frac{\bar{v}}{a} \frac{\partial \bar{\theta}}{\partial \phi} + \bar{w} \frac{\partial \bar{\theta}}{\partial z^*} - \bar{Q} = -\frac{1}{a \cos \phi} \overline{u' \frac{\partial \theta'}{\partial \lambda}} - \frac{1}{a} \overline{v' \frac{\partial \theta'}{\partial \phi}} - \overline{w' \frac{\partial \theta'}{\partial z^*}} \quad (\text{A.6e})$$

となる. ここで (A.5), (A.6) から東西平均からのずれに関する連続の式

$$\frac{1}{a \cos \phi} \left[\frac{\partial u'}{\partial \lambda} + \frac{\partial}{\partial \phi} (v' \cos \phi) \right] + \frac{1}{\rho_0} \frac{\partial}{\partial z^*} (\rho_0 w') = 0 \quad (\text{A.7})$$

が得られる.

(A.7) を使って (A.6) を変形する. (A.7) に u' をかけてオイラー平均をとると

$$\frac{1}{a \cos \phi} \overline{u' \frac{\partial u'}{\partial \lambda}} + \frac{1}{a} \overline{u' \frac{\partial v'}{\partial \phi}} - \frac{\tan \phi}{a} \overline{u' v'} + \overline{u' \frac{\partial w'}{\partial z^*}} + \frac{1}{\rho_0} \frac{\partial \rho_0}{\partial z^*} \overline{u' w'} = 0 \quad (\text{A.8})$$

これを (A.6) に加えると

$$\begin{aligned} \frac{\partial \bar{u}}{\partial t} + \frac{1}{a} \bar{v} \frac{\partial \bar{u}}{\partial \phi} + \bar{w} \frac{\partial \bar{u}}{\partial z^*} - f \bar{v} - \frac{\tan \phi}{a} \bar{u} \bar{v} - \bar{X} \\ = -\frac{2}{a \cos \phi} \overline{u' \frac{\partial u'}{\partial \lambda}} - \frac{1}{a} \overline{v' \frac{\partial u'}{\partial \phi}} - \overline{w' \frac{\partial u'}{\partial z^*}} - \frac{1}{a} \overline{u' \frac{\partial v'}{\partial \phi}} + \frac{2 \tan \phi}{a} \overline{u' v'} - \overline{u' \frac{\partial w'}{\partial z^*}} - \frac{1}{\rho_0} \frac{\partial \rho_0}{\partial z^*} \overline{u' w'} \end{aligned}$$

ここで

$$\begin{aligned} -\frac{2}{a \cos \phi} \overline{u' \frac{\partial u'}{\partial \lambda}} &= -\frac{1}{a \cos \phi} \frac{\partial (\overline{u'^2})}{\partial \lambda} = 0, \\ -\frac{1}{a} \overline{v' \frac{\partial u'}{\partial \phi}} - \frac{1}{a} \overline{u' \frac{\partial v'}{\partial \phi}} + \frac{2 \tan \phi}{a} \overline{u' v'} &= -\frac{1}{a \cos^2 \phi} \frac{\partial}{\partial \phi} (\overline{v' u'} \cos^2 \phi), \\ -\overline{w' \frac{\partial u'}{\partial z^*}} - \overline{u' \frac{\partial w'}{\partial z^*}} - \frac{1}{\rho_0} \frac{\partial \rho_0}{\partial z^*} \overline{u' w'} &= -\frac{1}{\rho_0} \frac{\partial}{\partial z^*} (\rho_0 \overline{w' u'}) \end{aligned}$$

を用いると,

$$\begin{aligned} \frac{\partial \bar{u}}{\partial t} + \frac{1}{a} \bar{v} \frac{\partial \bar{u}}{\partial \phi} + \bar{w} \frac{\partial \bar{u}}{\partial z^*} - f \bar{v} - \frac{\tan \phi}{a} \bar{u} \bar{v} - \bar{X} \\ = -\frac{1}{a \cos^2 \phi} \frac{\partial}{\partial \phi} (\overline{v'u'} \cos^2 \phi) - \frac{1}{\rho_0} \frac{\partial}{\partial z^*} (\rho_0 \overline{w'u'}) \end{aligned}$$

と書くことができる. (A.6) に関しても同様に, (A.7) に v' をかけてオイラー平均をとった式

$$\frac{1}{a \cos \phi} \overline{v' \frac{\partial u'}{\partial \lambda}} + \frac{1}{a} \overline{v' \frac{\partial v'}{\partial \phi}} + \frac{\tan \phi}{a} \overline{v'^2} + \overline{v' \frac{\partial w'}{\partial z^*}} + \frac{1}{\rho_0} \frac{\partial \rho_0}{\partial z^*} \overline{v'w'} = 0 \quad (\text{A.9})$$

を (A.6) に加えると

$$\begin{aligned} \frac{\partial \bar{v}}{\partial t} + \frac{\bar{v}}{a} \frac{\partial \bar{v}}{\partial \phi} + \bar{w} \frac{\partial \bar{v}}{\partial z^*} + f \bar{u} + \frac{\tan \phi}{a} (\bar{u})^2 + \frac{1}{a} \frac{\partial \bar{\Phi}}{\partial \phi} - \bar{Y} \\ = -\frac{1}{a \cos \phi} \overline{u' \frac{\partial v'}{\partial \lambda}} - \frac{1}{a} \overline{v' \frac{\partial v'}{\partial \phi}} - \overline{w' \frac{\partial v'}{\partial z^*}} - \frac{\tan \phi}{a} \overline{u'^2} \\ - \frac{1}{a \cos \phi} \overline{v' \frac{\partial u'}{\partial \lambda}} - \frac{1}{a} \overline{v' \frac{\partial v'}{\partial \phi}} + \frac{\tan \phi}{a} \overline{v'^2} - \overline{v' \frac{\partial w'}{\partial z^*}} - \frac{1}{\rho_0} \frac{\partial \rho_0}{\partial z^*} \overline{v'w'} \end{aligned}$$

が得られる. ここで

$$\begin{aligned} -\frac{1}{a \cos \phi} \overline{u' \frac{\partial v'}{\partial \lambda}} - \frac{1}{a \cos \phi} \overline{v' \frac{\partial u'}{\partial \lambda}} &= -\frac{1}{a \cos \phi} \frac{\partial (\overline{u'v'})}{\partial \lambda} = 0, \\ -\frac{1}{a} \overline{v' \frac{\partial v'}{\partial \phi}} - \frac{1}{a} \overline{v' \frac{\partial v'}{\partial \phi}} + \frac{\tan \phi}{a} \overline{v'^2} &= -\frac{1}{a \cos \phi} \frac{\partial}{\partial \phi} (\cos \phi \overline{v'^2}) \\ -\overline{w' \frac{\partial v'}{\partial z^*}} - \overline{v' \frac{\partial w'}{\partial z^*}} - \frac{1}{\rho_0} \frac{\partial \rho_0}{\partial z^*} \overline{v'w'} &= -\frac{1}{\rho_0} \frac{\partial}{\partial z^*} (\rho_0 \overline{v'w'}) \end{aligned} \quad (\text{A.10})$$

を用いると

$$\begin{aligned} \frac{\partial \bar{v}}{\partial t} + \frac{\bar{v}}{a} \frac{\partial \bar{v}}{\partial \phi} + \bar{w} \frac{\partial \bar{v}}{\partial z^*} + f \bar{u} + \frac{\tan \phi}{a} (\bar{u})^2 + \frac{1}{a} \frac{\partial \bar{\Phi}}{\partial \phi} - \bar{Y} \\ = -\frac{1}{a \cos \phi} \frac{\partial}{\partial \phi} (\cos \phi \overline{v'^2}) - \frac{\tan \phi}{a} \overline{u'^2} - \frac{1}{\rho_0} \frac{\partial}{\partial z^*} (\rho_0 \overline{v'w'}) \end{aligned}$$

と書くことができる. (A.6) についても同様に, (A.7) に θ' をかけてオイラー平均をとった式

$$\frac{1}{a \cos \phi} \overline{\theta' \frac{\partial u'}{\partial \lambda}} + \frac{1}{a} \overline{\theta' \frac{\partial v'}{\partial \phi}} - \frac{\tan \phi}{a} \overline{\theta'v'} + \overline{\theta' \frac{\partial w'}{\partial z^*}} + \frac{1}{\rho_0} \frac{\partial \rho_0}{\partial z^*} \overline{\theta'w'} = 0 \quad (\text{A.11})$$

を (A.6) に加えると

$$\begin{aligned} \frac{\partial \bar{\theta}}{\partial t} + \frac{\bar{v}}{a} \frac{\partial \bar{\theta}}{\partial \phi} + \bar{w} \frac{\partial \bar{\theta}}{\partial z^*} - \bar{Q} \\ = -\frac{1}{a \cos \phi} \overline{u' \frac{\partial \theta'}{\partial \lambda}} - \frac{1}{a} \overline{v' \frac{\partial \theta'}{\partial \phi}} - \overline{w' \frac{\partial \theta'}{\partial z^*}} \\ - \frac{1}{a \cos \phi} \overline{\theta' \frac{\partial u'}{\partial \lambda}} - \frac{1}{a} \overline{\theta' \frac{\partial v'}{\partial \phi}} + \frac{\tan \phi}{a} \overline{\theta' v'} - \overline{\theta' \frac{\partial w'}{\partial z^*}} - \frac{1}{\rho_0} \frac{\partial \rho_0}{\partial z^*} \overline{\theta' w'} \end{aligned}$$

が得られる. ここで

$$\begin{aligned} -\frac{1}{a \cos \phi} \overline{u' \frac{\partial \theta'}{\partial \lambda}} - \frac{1}{a \cos \phi} \overline{\theta' \frac{\partial u'}{\partial \lambda}} &= -\frac{1}{a \cos \phi} \frac{\partial (\overline{u' \theta'})}{\partial \lambda} = 0, \\ -\frac{1}{a} \overline{v' \frac{\partial \theta'}{\partial \phi}} - \frac{1}{a} \overline{\theta' \frac{\partial v'}{\partial \phi}} + \frac{\tan \phi}{a} \overline{\theta' v'} &= -\frac{1}{a \cos \phi} \frac{\partial}{\partial \phi} (\cos \phi \overline{v' \theta'}) \\ -\overline{w' \frac{\partial \theta'}{\partial z^*}} - \overline{\theta' \frac{\partial w'}{\partial z^*}} - \frac{1}{\rho_0} \frac{\partial \rho_0}{\partial z^*} \overline{\theta' w'} &= -\frac{1}{\rho_0} \frac{\partial}{\partial z^*} (\rho_0 \overline{w' \theta'}) \end{aligned}$$

を用いると

$$\frac{\partial \bar{\theta}}{\partial t} + \frac{\bar{v}}{a} \frac{\partial \bar{\theta}}{\partial \phi} + \bar{w} \frac{\partial \bar{\theta}}{\partial z^*} - \bar{Q} = -\frac{1}{a \cos \phi} \frac{\partial}{\partial \phi} (\cos \phi \overline{v' \theta'}) - \frac{1}{\rho_0} \frac{\partial}{\partial z^*} (\rho_0 \overline{w' \theta'})$$

となる.

以上をまとめると, 以下のオイラー平均方程式が得られる.

$$\begin{aligned} \frac{\partial \bar{u}}{\partial t} + \frac{1}{a} \bar{v} \frac{\partial \bar{u}}{\partial \phi} + \bar{w} \frac{\partial \bar{u}}{\partial z^*} - f \bar{v} - \frac{\tan \phi}{a} \bar{u} \bar{v} - \bar{X} \\ = -\frac{1}{a \cos^2 \phi} \frac{\partial}{\partial \phi} (\overline{v'u'} \cos^2 \phi) - \frac{1}{\rho_0} \frac{\partial}{\partial z^*} (\rho_0 \overline{w'u'}), \end{aligned} \quad (\text{A.12a})$$

$$\begin{aligned} \frac{\partial \bar{v}}{\partial t} + \frac{\bar{v}}{a} \frac{\partial \bar{v}}{\partial \phi} + \bar{w} \frac{\partial \bar{v}}{\partial z^*} + f \bar{u} + \frac{\tan \phi}{a} (\bar{u})^2 + \frac{1}{a} \frac{\partial \bar{\Phi}}{\partial \phi} - \bar{Y} \\ = -\frac{1}{a \cos \phi} \frac{\partial}{\partial \phi} (\overline{v'^2} \cos \phi) - \frac{1}{\rho_0} \frac{\partial}{\partial z^*} (\rho_0 \overline{v'w'}) - \frac{\overline{u'^2} \tan \phi}{a}, \end{aligned} \quad (\text{A.12b})$$

$$\frac{\partial \bar{\Phi}}{\partial z^*} - \frac{Re^{-\kappa z^*/H}}{H} \bar{\theta} = 0, \quad (\text{A.12c})$$

$$\frac{1}{a \cos \phi} \frac{\partial}{\partial \phi} (\bar{v} \cos \phi) + \frac{1}{\rho_0} \frac{\partial}{\partial z^*} (\rho_0 \bar{w}) = 0, \quad (\text{A.12d})$$

$$\frac{\partial \bar{\theta}}{\partial t} + \frac{\bar{v}}{a} \frac{\partial \bar{\theta}}{\partial \phi} + \bar{w} \frac{\partial \bar{\theta}}{\partial z^*} - \bar{Q} = -\frac{1}{a \cos \phi} \frac{\partial}{\partial \phi} (\overline{v'\theta'}) \cos \phi - \frac{1}{\rho_0} \frac{\partial}{\partial z^*} (\rho_0 \overline{w'\theta'}). \quad (\text{A.12e})$$

1.3 変形オイラー平均方程式系

(A.12) を EP フラックス, 残差循環を用いて書き直す. EP フラックス, 残差循環は以下のように定義する.

$$\bar{v}^* = \bar{v} - \frac{1}{\rho_0} \frac{\partial}{\partial z^*} \left(\rho_0 \frac{\overline{v'\theta'}}{\frac{\partial \theta}{\partial z^*}} \right) \quad (\text{A.13a})$$

$$\bar{w}^* = \bar{w} + \frac{1}{a \cos \phi} \frac{\partial}{\partial \phi} \left(\cos \phi \frac{\overline{v'\theta'}}{\frac{\partial \theta}{\partial z^*}} \right) \quad (\text{A.13b})$$

$$\begin{aligned} F_\phi &= \rho_0 a \cos \phi \left(\frac{\partial \bar{u}}{\partial z^*} \frac{\overline{v'\theta'}}{\frac{\partial \theta}{\partial z^*}} - \overline{u'v'} \right) \\ F_z^* &= \rho_0 a \cos \phi \left(\left[f - \frac{\frac{\partial \bar{u} \cos \phi}{\partial \phi}}{a \cos \phi} \right] \frac{\overline{v'\theta'}}{\frac{\partial \theta}{\partial z^*}} - \overline{u'w'} \right) \end{aligned}$$

まず連続の式を書き換える. (A.12) に (A.13), (A.13) を代入すると

$$\begin{aligned} & \frac{1}{a \cos \phi} \frac{\partial}{\partial \phi} \left[\left\{ \bar{v}^* + \frac{1}{\rho_0} \frac{\partial}{\partial z^*} \left(\rho_0 \frac{\overline{v'\theta'}}{\frac{\partial \theta}{\partial z^*}} \right) \right\} \cos \phi \right] \\ & + \frac{1}{\rho_0} \frac{\partial}{\partial z^*} \left[\rho_0 \left\{ \bar{w}^* - \frac{1}{a \cos \phi} \frac{\partial}{\partial \phi} \left(\cos \phi \frac{\overline{v'\theta'}}{\frac{\partial \theta}{\partial z^*}} \right) \right\} \right] = 0, \\ & \frac{1}{a \cos \phi} \frac{\partial}{\partial \phi} (\bar{v}^* \cos \phi) + \frac{1}{\rho_0} \frac{\partial}{\partial z^*} (\rho_0 \bar{w}^*) \\ & + \frac{1}{a \cos \phi} \frac{\partial}{\partial \phi} \left\{ \frac{1}{\rho_0} \frac{\partial}{\partial z^*} \left(\rho_0 \frac{\overline{v'\theta'}}{\frac{\partial \theta}{\partial z^*}} \right) \cos \phi \right\} - \frac{1}{\rho_0} \frac{\partial}{\partial z^*} \left\{ \rho_0 \frac{1}{a \cos \phi} \frac{\partial}{\partial \phi} \left(\cos \phi \frac{\overline{v'\theta'}}{\frac{\partial \theta}{\partial z^*}} \right) \right\} = 0. \end{aligned}$$

この第三項と第四項だけを取り出すと

$$\begin{aligned} & \frac{1}{a \cos \phi} \frac{\partial}{\partial \phi} \left\{ \frac{1}{\rho_0} \frac{\partial}{\partial z^*} \left(\rho_0 \frac{\overline{v'\theta'}}{\frac{\partial \theta}{\partial z^*}} \right) \cos \phi \right\} - \frac{1}{\rho_0} \frac{\partial}{\partial z^*} \left\{ \rho_0 \frac{1}{a \cos \phi} \frac{\partial}{\partial \phi} \left(\cos \phi \frac{\overline{v'\theta'}}{\frac{\partial \theta}{\partial z^*}} \right) \right\} \\ & = \frac{1}{a \cos \phi} \left[\frac{\partial}{\partial \phi} \left\{ \frac{1}{\rho_0} \frac{\partial}{\partial z^*} \left(\rho_0 \frac{\overline{v'\theta'}}{\frac{\partial \theta}{\partial z^*}} \right) \cos \phi \right\} - \frac{1}{\rho_0} \frac{\partial}{\partial z^*} \left\{ \rho_0 \frac{\partial}{\partial \phi} \left(\cos \phi \frac{\overline{v'\theta'}}{\frac{\partial \theta}{\partial z^*}} \right) \right\} \right] \\ & = \frac{1}{a \cos \phi} \left[\frac{1}{\rho_0} \frac{\partial}{\partial \phi} \left\{ \frac{\partial}{\partial z^*} \left(\rho_0 \frac{\overline{v'\theta'}}{\frac{\partial \theta}{\partial z^*}} \cos \phi \right) \right\} - \frac{1}{\rho_0} \frac{\partial}{\partial z^*} \left\{ \frac{\partial}{\partial \phi} \left(\rho_0 \cos \phi \frac{\overline{v'\theta'}}{\frac{\partial \theta}{\partial z^*}} \right) \right\} \right] \\ & = 0. \end{aligned}$$

したがって, 連続の式は以下のようになる.

$$\frac{1}{a \cos \phi} \frac{\partial}{\partial \phi} (\bar{v}^* \cos \phi) + \frac{1}{\rho_0} \frac{\partial}{\partial z^*} (\rho_0 \bar{w}^*) = 0. \quad (\text{A.14})$$

次に u の式を書き換える. (A.12) に (A.13), (A.13) を代入すると

$$\begin{aligned}
& \frac{\partial \bar{u}}{\partial t} + \frac{1}{a} \left[\bar{v}^* + \frac{1}{\rho_0} \frac{\partial}{\partial z^*} \left(\rho_0 \frac{\overline{v'\theta'}}{\frac{\partial \theta}{\partial z^*}} \right) \right] \frac{\partial \bar{u}}{\partial \phi} + \left[\bar{w}^* - \frac{1}{a \cos \phi} \frac{\partial}{\partial \phi} \left(\cos \phi \frac{\overline{v'\theta'}}{\frac{\partial \theta}{\partial z^*}} \right) \right] \frac{\partial \bar{u}}{\partial z^*} \\
& \quad - f \left[\bar{v}^* + \frac{1}{\rho_0} \frac{\partial}{\partial z^*} \left(\rho_0 \frac{\overline{v'\theta'}}{\frac{\partial \theta}{\partial z^*}} \right) \right] - \frac{\tan \phi}{a} \bar{u} \left[\bar{v}^* + \frac{1}{\rho_0} \frac{\partial}{\partial z^*} \left(\rho_0 \frac{\overline{v'\theta'}}{\frac{\partial \theta}{\partial z^*}} \right) \right] - \bar{X} \\
& = -\frac{1}{a \cos^2 \phi} \frac{\partial}{\partial \phi} (\overline{v'u'} \cos^2 \phi) - \frac{1}{\rho_0} \frac{\partial}{\partial z^*} (\rho_0 \overline{w'u'}), \\
& \frac{\partial \bar{u}}{\partial t} + \frac{\bar{v}^*}{a} \frac{\partial \bar{u}}{\partial \phi} + \bar{w}^* \frac{\partial \bar{u}}{\partial z^*} - f \bar{v}^* - \frac{\tan \phi}{a} \bar{u} \bar{v}^* - \bar{X} \\
& = -\frac{1}{a \cos^2 \phi} \frac{\partial}{\partial \phi} (\overline{v'u'} \cos^2 \phi) + \frac{1}{a \cos \phi} \frac{\partial}{\partial \phi} \left(\cos \phi \frac{\overline{v'\theta'}}{\frac{\partial \theta}{\partial z^*}} \right) \frac{\partial \bar{u}}{\partial z^*} \\
& \quad + f \frac{1}{\rho_0} \frac{\partial}{\partial z^*} \left(\rho_0 \frac{\overline{v'\theta'}}{\frac{\partial \theta}{\partial z^*}} \right) - \frac{1}{\rho_0} \frac{\partial}{\partial z^*} (\rho_0 \overline{w'u'}) \\
& \quad - \frac{1}{\rho_0 a} \frac{\partial}{\partial z^*} \left(\rho_0 \frac{\overline{v'\theta'}}{\frac{\partial \theta}{\partial z^*}} \right) \frac{\partial \bar{u}}{\partial \phi} + \frac{\tan \phi}{a} \bar{u} \frac{1}{\rho_0} \frac{\partial}{\partial z^*} \left(\rho_0 \frac{\overline{v'\theta'}}{\frac{\partial \theta}{\partial z^*}} \right), \\
& \frac{\partial \bar{u}}{\partial t} + \frac{\bar{v}^*}{a \cos \phi} \frac{\partial}{\partial \phi} (\bar{u} \cos \phi) + \bar{w}^* \frac{\partial \bar{u}}{\partial z^*} - f \bar{v}^* - \bar{X} \\
& = -\frac{1}{\rho_0 a^2 \cos^2 \phi} \frac{\partial}{\partial \phi} (\rho_0 a \overline{v'u'} \cos^2 \phi) + \frac{1}{a \cos \phi} \frac{\partial}{\partial \phi} \left(\cos \phi \frac{\overline{v'\theta'}}{\frac{\partial \theta}{\partial z^*}} \right) \frac{\partial \bar{u}}{\partial z^*} \\
& \quad + \frac{1}{\rho_0 a \cos \phi} \frac{\partial}{\partial z^*} \left(f \rho_0 a \cos \phi \frac{\overline{v'\theta'}}{\frac{\partial \theta}{\partial z^*}} \right) - \frac{1}{\rho_0 a \cos \phi} \frac{\partial}{\partial z^*} (\rho_0 a \cos \phi \overline{w'u'}) \\
& \quad - \frac{1}{\rho_0 a} \frac{\partial}{\partial z^*} \left(\rho_0 \frac{\overline{v'\theta'}}{\frac{\partial \theta}{\partial z^*}} \right) \frac{\partial \bar{u}}{\partial \phi} + \frac{\tan \phi}{a} \bar{u} \frac{1}{\rho_0} \frac{\partial}{\partial z^*} \left(\rho_0 \frac{\overline{v'\theta'}}{\frac{\partial \theta}{\partial z^*}} \right) \tag{A.15}
\end{aligned}$$

(1.3) の右辺を以下のように変形する.

$$\begin{aligned}
& -\frac{1}{\rho_0 a^2 \cos^2 \phi} \frac{\partial}{\partial \phi} (\rho_0 a \overline{v' u'} \cos^2 \phi) + \frac{1}{\rho_0 a^2 \cos^2 \phi} \rho_0 a \cos \phi \frac{\partial \bar{u}}{\partial z^*} \frac{\partial}{\partial \phi} \left(\cos \phi \frac{\overline{v' \theta'}}{\frac{\partial \theta}{\partial z^*}} \right) \\
& \quad + \frac{1}{\rho_0 a \cos \phi} \frac{\partial}{\partial z^*} \left(f \rho_0 a \cos \phi \frac{\overline{v' \theta'}}{\frac{\partial \theta}{\partial z^*}} \right) - \frac{1}{\rho_0 a \cos \phi} \frac{\partial}{\partial z^*} (\rho_0 a \cos \phi \overline{w' u'}) \\
& \quad - \frac{1}{\rho_0 a} \frac{\partial}{\partial z^*} \left(\rho_0 \frac{\overline{v' \theta'}}{\frac{\partial \theta}{\partial z^*}} \frac{\partial \bar{u}}{\partial \phi} \right) + \frac{1}{\rho_0 a} \rho_0 \frac{\overline{v' \theta'}}{\frac{\partial \theta}{\partial z^*}} \frac{\partial}{\partial z^*} \left(\frac{\partial \bar{u}}{\partial \phi} \right) \\
& \quad + \frac{\tan \phi}{\rho_0 a} \frac{\partial}{\partial z^*} \left(\bar{u} \rho_0 \frac{\overline{v' \theta'}}{\frac{\partial \theta}{\partial z^*}} \right) - \frac{\tan \phi}{\rho_0 a} \rho_0 \frac{\overline{v' \theta'}}{\frac{\partial \theta}{\partial z^*}} \frac{\partial}{\partial z^*} (\bar{u}) \\
& = \frac{1}{\rho_0 a^2 \cos^2 \phi} \left[-\frac{\partial}{\partial \phi} (\rho_0 a \overline{v' u'} \cos^2 \phi) + \rho_0 a \cos \phi \frac{\partial \bar{u}}{\partial z^*} \frac{\partial}{\partial \phi} \left(\cos \phi \frac{\overline{v' \theta'}}{\frac{\partial \theta}{\partial z^*}} \right) \right] \\
& \quad + \frac{1}{\rho_0 a} \rho_0 \frac{\overline{v' \theta'}}{\frac{\partial \theta}{\partial z^*}} \frac{\partial}{\partial z^*} \left(\frac{\partial \bar{u}}{\partial \phi} \right) - \frac{\tan \phi}{\rho_0 a} \rho_0 \frac{\overline{v' \theta'}}{\frac{\partial \theta}{\partial z^*}} \frac{\partial \bar{u}}{\partial z^*} \\
& \quad + \frac{1}{\rho_0 a \cos \phi} \frac{\partial}{\partial z^*} \left[\left(f \rho_0 a \cos \phi \frac{\overline{v' \theta'}}{\frac{\partial \theta}{\partial z^*}} \right) - \rho_0 a \cos \phi \overline{w' u'} \right] \\
& \quad - \frac{1}{\rho_0 a} \frac{\partial}{\partial z^*} \left(\rho_0 \frac{\overline{v' \theta'}}{\frac{\partial \theta}{\partial z^*}} \frac{\partial \bar{u}}{\partial \phi} \right) + \frac{\tan \phi}{\rho_0 a} \frac{\partial}{\partial z^*} \left(\bar{u} \rho_0 \frac{\overline{v' \theta'}}{\frac{\partial \theta}{\partial z^*}} \right) \\
& = \frac{1}{\rho_0 a^2 \cos^2 \phi} \left[-\frac{\partial}{\partial \phi} (\rho_0 a \overline{v' u'} \cos^2 \phi) + \rho_0 a \cos \phi \frac{\partial \bar{u}}{\partial z^*} \frac{\partial}{\partial \phi} \left(\cos \phi \frac{\overline{v' \theta'}}{\frac{\partial \theta}{\partial z^*}} \right) \right] \\
& \quad + \frac{1}{\rho_0 a^2 \cos^2 \phi} \left[\rho_0 a \cos^2 \phi \frac{\overline{v' \theta'}}{\frac{\partial \theta}{\partial z^*}} \frac{\partial}{\partial z^*} \left(\frac{\partial \bar{u}}{\partial \phi} \right) - \rho_0 a \cos^2 \phi \tan \phi \frac{\overline{v' \theta'}}{\frac{\partial \theta}{\partial z^*}} \frac{\partial \bar{u}}{\partial z^*} \right] \\
& \quad + \frac{1}{\rho_0 a \cos \phi} \frac{\partial}{\partial z^*} \left[\left(f \rho_0 a \cos \phi \frac{\overline{v' \theta'}}{\frac{\partial \theta}{\partial z^*}} \right) - \rho_0 a \cos \phi \overline{w' u'} \right] \\
& \quad + \frac{1}{\rho_0 a \cos \phi} \left[-\cos \phi \frac{\partial}{\partial z^*} \left(\rho_0 \frac{\overline{v' \theta'}}{\frac{\partial \theta}{\partial z^*}} \frac{\partial \bar{u}}{\partial \phi} \right) + \cos \phi \tan \phi \frac{\partial}{\partial z^*} \left(\bar{u} \rho_0 \frac{\overline{v' \theta'}}{\frac{\partial \theta}{\partial z^*}} \right) \right] \\
& = \frac{1}{\rho_0 a^2 \cos^2 \phi} \left[-\frac{\partial}{\partial \phi} (\rho_0 a \overline{v' u'} \cos^2 \phi) + \rho_0 a \cos \phi \frac{\partial \bar{u}}{\partial z^*} \frac{\partial}{\partial \phi} \left(\cos \phi \frac{\overline{v' \theta'}}{\frac{\partial \theta}{\partial z^*}} \right) \right] \\
& \quad + \frac{1}{\rho_0 a^2 \cos^2 \phi} \left[\rho_0 a \cos^2 \phi \frac{\overline{v' \theta'}}{\frac{\partial \theta}{\partial z^*}} \frac{\partial}{\partial \phi} \left(\frac{\partial \bar{u}}{\partial z^*} \right) + \cos \phi \frac{\partial}{\partial \phi} (\rho_0 a \cos \phi) \frac{\overline{v' \theta'}}{\frac{\partial \theta}{\partial z^*}} \frac{\partial \bar{u}}{\partial z^*} \right] \\
& \quad + \frac{1}{\rho_0 a \cos \phi} \frac{\partial}{\partial z^*} \left[f \rho_0 a \cos \phi \frac{\overline{v' \theta'}}{\frac{\partial \theta}{\partial z^*}} - \rho_0 a \cos \phi \overline{w' u'} \right] \\
& \quad + \frac{1}{\rho_0 a \cos \phi} \frac{\partial}{\partial z^*} \left[-\rho_0 \cos \phi \frac{\overline{v' \theta'}}{\frac{\partial \theta}{\partial z^*}} \frac{\partial \bar{u}}{\partial \phi} + \sin \phi \bar{u} \rho_0 \frac{\overline{v' \theta'}}{\frac{\partial \theta}{\partial z^*}} \right] \tag{A.16}
\end{aligned}$$

(1.3) の第一項と第二項だけ取り出すと

$$\begin{aligned}
& \frac{1}{\rho_0 a^2 \cos^2 \phi} \left[-\frac{\partial}{\partial \phi} (\rho_0 a \overline{v' u'}) \cos^2 \phi + \rho_0 a \cos \phi \frac{\partial \bar{u}}{\partial z^*} \frac{\partial}{\partial \phi} \left(\cos \phi \frac{\overline{v' \theta'}}{\frac{\partial \theta}{\partial z^*}} \right) \right] \\
& + \frac{1}{\rho_0 a^2 \cos^2 \phi} \left[\rho_0 a \cos^2 \phi \frac{\overline{v' \theta'}}{\frac{\partial \theta}{\partial z^*}} \frac{\partial}{\partial \phi} \left(\frac{\partial \bar{u}}{\partial z^*} \right) + \cos \phi \frac{\partial}{\partial \phi} (\rho_0 a \cos \phi) \frac{\overline{v' \theta'}}{\frac{\partial \theta}{\partial z^*}} \frac{\partial \bar{u}}{\partial z^*} \right] \\
& = \frac{1}{\rho_0 a^2 \cos^2 \phi} \left[-\frac{\partial}{\partial \phi} (\rho_0 a \overline{v' u'}) \cos^2 \phi \right] \\
& + \frac{1}{\rho_0 a^2 \cos^2 \phi} \left[\rho_0 a \cos^2 \phi \frac{\overline{v' \theta'}}{\frac{\partial \theta}{\partial z^*}} \frac{\partial}{\partial \phi} \left(\frac{\partial \bar{u}}{\partial z^*} \right) + \frac{\partial \bar{u}}{\partial z^*} \frac{\partial}{\partial \phi} \left(\rho_0 a \cos^2 \phi \frac{\overline{v' \theta'}}{\frac{\partial \theta}{\partial z^*}} \right) \right] \\
& = \frac{1}{\rho_0 a^2 \cos^2 \phi} \left[-\frac{\partial}{\partial \phi} (\rho_0 a \overline{v' u'}) \cos^2 \phi \right] + \frac{1}{\rho_0 a^2 \cos^2 \phi} \left[\frac{\partial}{\partial \phi} \left(\rho_0 a \cos^2 \phi \frac{\overline{v' \theta'}}{\frac{\partial \theta}{\partial z^*}} \frac{\partial \bar{u}}{\partial z^*} \right) \right] \\
& = \frac{1}{\rho_0 a^2 \cos^2 \phi} \frac{\partial}{\partial \phi} \left[-\rho_0 a \overline{v' u'} \cos^2 \phi + \rho_0 a \cos^2 \phi \frac{\overline{v' \theta'}}{\frac{\partial \theta}{\partial z^*}} \frac{\partial \bar{u}}{\partial z^*} \right] \\
& = \frac{1}{\rho_0 a^2 \cos^2 \phi} \frac{\partial}{\partial \phi} \left[\rho_0 a \cos^2 \phi \left\{ \frac{\partial \bar{u}}{\partial z^*} \frac{\overline{v' \theta'}}{\frac{\partial \theta}{\partial z^*}} - \overline{v' u'} \right\} \right] \\
& = \frac{1}{\rho_0 a^2 \cos^2 \phi} \frac{\partial}{\partial \phi} (\cos \phi F_\phi^*)
\end{aligned}$$

(1.3) の第三項と第四項だけ取り出すと

$$\begin{aligned}
& \frac{1}{\rho_0 a \cos \phi} \frac{\partial}{\partial z^*} \left[f \rho_0 a \cos \phi \frac{\overline{v' \theta'}}{\frac{\partial \theta}{\partial z^*}} - \rho_0 a \cos \phi \overline{w' u'} \right] + \frac{1}{\rho_0 a \cos \phi} \frac{\partial}{\partial z^*} \left[-\rho_0 \cos \phi \frac{\overline{v' \theta'}}{\frac{\partial \theta}{\partial z^*}} \frac{\partial \bar{u}}{\partial \phi} + \sin \phi \bar{u} \rho_0 \frac{\overline{v' \theta'}}{\frac{\partial \theta}{\partial z^*}} \right] \\
& = \frac{1}{\rho_0 a \cos \phi} \frac{\partial}{\partial z^*} \left[\rho_0 a \cos \phi \left\{ f \frac{\overline{v' \theta'}}{\frac{\partial \theta}{\partial z^*}} - \overline{w' u'} - \frac{\overline{v' \theta'}}{a \frac{\partial \theta}{\partial z^*}} \frac{\partial \bar{u}}{\partial \phi} + \sin \phi \bar{u} \frac{\overline{v' \theta'}}{a \cos \phi \frac{\partial \theta}{\partial z^*}} \right\} \right] \\
& = \frac{1}{\rho_0 a \cos \phi} \frac{\partial}{\partial z^*} \left[\rho_0 a \cos \phi \left\{ f \frac{\overline{v' \theta'}}{\frac{\partial \theta}{\partial z^*}} - \left(\cos \phi \frac{\partial \bar{u}}{\partial \phi} - \sin \phi \bar{u} \right) \frac{\overline{v' \theta'}}{a \cos \phi \frac{\partial \theta}{\partial z^*}} - \overline{w' u'} \right\} \right] \\
& = \frac{1}{\rho_0 a \cos \phi} \frac{\partial}{\partial z^*} \left[\rho_0 a \cos \phi \left\{ f \frac{\overline{v' \theta'}}{\frac{\partial \theta}{\partial z^*}} - \frac{\partial (\bar{u} \cos \phi)}{\partial \phi} \frac{\overline{v' \theta'}}{a \cos \phi \frac{\partial \theta}{\partial z^*}} - \overline{w' u'} \right\} \right] \\
& = \frac{1}{\rho_0 a \cos \phi} \frac{\partial}{\partial z^*} \left[\rho_0 a \cos \phi \left\{ \left(f - \frac{\partial (\bar{u} \cos \phi)}{\partial \phi} \right) \frac{\overline{v' \theta'}}{a \cos \phi} \frac{\partial \theta}{\partial z^*} - \overline{w' u'} \right\} \right] \\
& = \frac{1}{\rho_0 a \cos \phi} \frac{\partial F_z^*}{\partial z^*}
\end{aligned}$$

以上より, (1.3) は次のようになる.

$$\begin{aligned}\frac{\partial \bar{u}}{\partial t} + \frac{\bar{v}^*}{a \cos \phi} \frac{\partial}{\partial \phi} (\bar{u} \cos \phi) + \bar{w}^* \frac{\partial \bar{u}}{\partial z^*} - f \bar{v}^* - \bar{X} &= \frac{1}{\rho_0 a^2 \cos^2 \phi} \frac{\partial}{\partial \phi} (\cos \phi F_\phi^*) + \frac{1}{\rho_0 a \cos \phi} \frac{\partial F_z^*}{\partial z^*}, \\ \frac{\partial \bar{u}}{\partial t} + \frac{\bar{v}^*}{a \cos \phi} \frac{\partial}{\partial \phi} (\bar{u} \cos \phi) + \bar{w}^* \frac{\partial \bar{u}}{\partial z^*} - f \bar{v}^* - \bar{X} &= \frac{1}{\rho_0 a \cos \phi} \nabla \cdot \mathbf{F}.\end{aligned}$$

ここで, 子午面内の発散を以下のように表した.

$$\nabla \cdot \mathbf{F} = \frac{1}{a \cos \phi} \frac{\partial (\cos \phi F_\phi)}{\partial \phi} + \frac{\partial F_z}{\partial z} \quad (\text{A.17})$$

次に熱力学の式を書き換える. (A.12) に (A.13), (A.13) を代入すると

$$\begin{aligned}\frac{\partial \bar{\theta}}{\partial t} + \frac{1}{a} \left[\bar{v}^* + \frac{1}{\rho_0} \frac{\partial}{\partial z^*} \left(\rho_0 \frac{\overline{v'\theta'}}{\frac{\partial \theta}{\partial z^*}} \right) \right] \frac{\partial \bar{\theta}}{\partial \phi} + \left[\bar{w}^* - \frac{1}{a \cos \phi} \frac{\partial}{\partial \phi} \left(\cos \phi \frac{\overline{v'\theta'}}{\frac{\partial \theta}{\partial z^*}} \right) \right] \frac{\partial \bar{\theta}}{\partial z^*} - \bar{Q} \\ = -\frac{1}{a \cos \phi} \frac{\partial}{\partial \phi} (\overline{v'\theta'} \cos \phi) - \frac{1}{\rho_0} \frac{\partial}{\partial z^*} (\rho_0 \overline{w'\theta'}), \\ \frac{\partial \bar{\theta}}{\partial t} + \frac{\bar{v}^*}{a} \frac{\partial \bar{\theta}}{\partial \phi} + \bar{w}^* \frac{\partial \bar{\theta}}{\partial z^*} - \bar{Q} \\ = -\frac{1}{\rho_0 a} \frac{\partial}{\partial z^*} \left(\rho_0 \frac{\overline{v'\theta'}}{\frac{\partial \theta}{\partial z^*}} \right) \frac{\partial \bar{\theta}}{\partial \phi} + \frac{1}{a \cos \phi} \frac{\partial}{\partial \phi} \left(\cos \phi \frac{\overline{v'\theta'}}{\frac{\partial \theta}{\partial z^*}} \right) \frac{\partial \bar{\theta}}{\partial z^*} \\ - \frac{1}{a \cos \phi} \frac{\partial}{\partial \phi} (\overline{v'\theta'} \cos \phi) - \frac{1}{\rho_0} \frac{\partial}{\partial z^*} (\rho_0 \overline{w'\theta'})\end{aligned}$$

となる. この右辺を更に変形すると

$$\begin{aligned}
& -\frac{1}{\rho_0} \frac{\partial}{\partial z^*} \left(\rho_0 \frac{\overline{v'\theta'}}{a \frac{\partial \bar{\theta}}{\partial z^*}} \right) \frac{\partial \bar{\theta}}{\partial \phi} + \frac{1}{a \cos \phi} \frac{\partial}{\partial \phi} \left(\cos \phi \frac{\overline{v'\theta'}}{\frac{\partial \bar{\theta}}{\partial z^*}} \right) \frac{\partial \bar{\theta}}{\partial z^*} \\
& \quad - \frac{1}{a \cos \phi} \frac{\partial}{\partial \phi} (\overline{v'\theta'} \cos \phi) - \frac{1}{\rho_0} \frac{\partial}{\partial z^*} (\rho_0 \overline{w'\theta'}) \\
= & -\frac{1}{\rho_0} \frac{\partial}{\partial z^*} \left(\rho_0 \frac{\overline{v'\theta'}}{a \frac{\partial \bar{\theta}}{\partial z^*}} \frac{\partial \bar{\theta}}{\partial \phi} \right) + \frac{\overline{v'\theta'}}{a \frac{\partial \bar{\theta}}{\partial z^*}} \frac{\partial}{\partial z^*} \frac{\partial \bar{\theta}}{\partial \phi} \\
& \quad + \frac{1}{a \cos \phi} \left[\frac{\partial}{\partial \phi} (\cos \phi \overline{v'\theta'}) \frac{1}{\frac{\partial \bar{\theta}}{\partial z^*}} + \cos \phi \overline{v'\theta'} \frac{\partial}{\partial \phi} \left(\frac{\partial \bar{\theta}}{\partial z^*} \right)^{-1} \right] \frac{\partial \bar{\theta}}{\partial z^*} \\
& \quad - \frac{1}{a \cos \phi} \frac{\partial}{\partial \phi} (\overline{v'\theta'} \cos \phi) - \frac{1}{\rho_0} \frac{\partial}{\partial z^*} (\rho_0 \overline{w'\theta'}) \\
= & -\frac{1}{\rho_0} \frac{\partial}{\partial z^*} \left(\rho_0 \frac{\overline{v'\theta'}}{a \frac{\partial \bar{\theta}}{\partial z^*}} \frac{\partial \bar{\theta}}{\partial \phi} \right) + \frac{\overline{v'\theta'}}{a \frac{\partial \bar{\theta}}{\partial z^*}} \frac{\partial}{\partial z^*} \frac{\partial \bar{\theta}}{\partial \phi} + \frac{1}{a} \overline{v'\theta'} \frac{\partial}{\partial \phi} \left(\frac{\partial \bar{\theta}}{\partial z^*} \right)^{-1} \frac{\partial \bar{\theta}}{\partial z^*} - \frac{1}{\rho_0} \frac{\partial}{\partial z^*} (\rho_0 \overline{w'\theta'}) \\
= & -\frac{1}{\rho_0} \frac{\partial}{\partial z^*} \left[\rho_0 \frac{\overline{v'\theta'}}{a \frac{\partial \bar{\theta}}{\partial z^*}} \frac{\partial \bar{\theta}}{\partial \phi} + \rho_0 \overline{w'\theta'} \right] + \frac{\overline{v'\theta'}}{a} \left[\frac{1}{\frac{\partial \bar{\theta}}{\partial z^*}} \frac{\partial}{\partial z^*} \frac{\partial \bar{\theta}}{\partial \phi} + \frac{\partial}{\partial \phi} \left(\frac{\partial \bar{\theta}}{\partial z^*} \right)^{-1} \frac{\partial \bar{\theta}}{\partial z^*} \right] \\
= & -\frac{1}{\rho_0} \frac{\partial}{\partial z^*} \left[\rho_0 \left(\frac{\overline{v'\theta'}}{a \frac{\partial \bar{\theta}}{\partial z^*}} \frac{\partial \bar{\theta}}{\partial \phi} + \overline{w'\theta'} \right) \right] + \frac{\overline{v'\theta'}}{a} \frac{\partial}{\partial \phi} \left(\frac{\partial \bar{\theta}}{\partial z^*} \right) \\
= & -\frac{1}{\rho_0} \frac{\partial}{\partial z^*} \left[\rho_0 \left(\frac{\overline{v'\theta'}}{a \frac{\partial \bar{\theta}}{\partial z^*}} \frac{\partial \bar{\theta}}{\partial \phi} + \overline{w'\theta'} \right) \right].
\end{aligned}$$

これより, 熱力学の式は以下のようなになる.

$$\frac{\partial \bar{\theta}}{\partial t} + \bar{v}^* \frac{\partial \bar{\theta}}{\partial \phi} + \bar{w}^* \frac{\partial \bar{\theta}}{\partial z^*} - \bar{Q} = -\frac{1}{\rho_0} \frac{\partial}{\partial z^*} \left[\rho_0 \left(\frac{\overline{v'\theta'}}{a \frac{\partial \bar{\theta}}{\partial z^*}} \frac{\partial \bar{\theta}}{\partial \phi} + \overline{w'\theta'} \right) \right].$$

最後に v の式について考える. (A.12) に (A.13), (A.13) を代入すると

$$\begin{aligned}
& \frac{\partial}{\partial t} \left[\bar{v}^* + \frac{1}{\rho_0} \frac{\partial}{\partial z^*} \left(\rho_0 \frac{\overline{v'\theta'}}{\frac{\partial\theta}{\partial z^*}} \right) \right] + \frac{1}{a} \left[\bar{v}^* + \frac{1}{\rho_0} \frac{\partial}{\partial z^*} \left(\rho_0 \frac{\overline{v'\theta'}}{\frac{\partial\theta}{\partial z^*}} \right) \right] \frac{\partial}{\partial\phi} \left[\bar{v}^* + \frac{1}{\rho_0} \frac{\partial}{\partial z^*} \left(\rho_0 \frac{\overline{v'\theta'}}{\frac{\partial\theta}{\partial z^*}} \right) \right] \\
& + \left[\bar{w}^* - \frac{1}{a \cos\phi} \frac{\partial}{\partial\phi} \left(\cos\phi \frac{\overline{v'\theta'}}{\frac{\partial\theta}{\partial z^*}} \right) \right] \frac{\partial}{\partial z^*} \left[\bar{v}^* + \frac{1}{\rho_0} \frac{\partial}{\partial z^*} \left(\rho_0 \frac{\overline{v'\theta'}}{\frac{\partial\theta}{\partial z^*}} \right) \right] \\
& + f\bar{u} + \frac{\tan\phi}{a} (\bar{u})^2 + \frac{1}{a} \frac{\partial\bar{\Phi}}{\partial\phi} - \bar{Y} \\
& = -\frac{1}{a \cos\phi} \frac{\partial}{\partial\phi} (\bar{v}^2 \cos\phi) - \frac{1}{\rho_0} \frac{\partial}{\partial z^*} (\rho_0 \overline{v'w'}) - \frac{\bar{u}^2 \tan\phi}{a}, \\
& f\bar{u} + \frac{\tan\phi}{a} (\bar{u})^2 + \frac{1}{a} \frac{\partial\bar{\Phi}}{\partial\phi} \\
& = -\frac{\partial}{\partial t} \left[\bar{v}^* + \frac{1}{\rho_0} \frac{\partial}{\partial z^*} \left(\rho_0 \frac{\overline{v'\theta'}}{\frac{\partial\theta}{\partial z^*}} \right) \right] - \frac{1}{a} \left[\bar{v}^* + \frac{1}{\rho_0} \frac{\partial}{\partial z^*} \left(\rho_0 \frac{\overline{v'\theta'}}{\frac{\partial\theta}{\partial z^*}} \right) \right] \frac{\partial}{\partial\phi} \left[\bar{v}^* + \frac{1}{\rho_0} \frac{\partial}{\partial z^*} \left(\rho_0 \frac{\overline{v'\theta'}}{\frac{\partial\theta}{\partial z^*}} \right) \right] \\
& - \left[\bar{w}^* - \frac{1}{a \cos\phi} \frac{\partial}{\partial\phi} \left(\cos\phi \frac{\overline{v'\theta'}}{\frac{\partial\theta}{\partial z^*}} \right) \right] \frac{\partial}{\partial z^*} \left[\bar{v}^* + \frac{1}{\rho_0} \frac{\partial}{\partial z^*} \left(\rho_0 \frac{\overline{v'\theta'}}{\frac{\partial\theta}{\partial z^*}} \right) \right] \\
& - \frac{1}{a \cos\phi} \frac{\partial}{\partial\phi} (\bar{v}^2 \cos\phi) - \frac{1}{\rho_0} \frac{\partial}{\partial z^*} (\rho_0 \overline{v'w'}) - \frac{\bar{u}^2 \tan\phi}{a} + \bar{Y}
\end{aligned}$$

Andrews *et al.* (1987) によれば, この式の右辺の量は左辺に比べれば小さい. 右辺の項を全てまとめて G と書くと v の式は次のようになる.

$$\bar{u} \left(f + \frac{\tan\phi}{a} \bar{u} \right) + \frac{1}{a} \frac{\partial\bar{\Phi}}{\partial\phi} = G.$$

以上をまとめると, 以下の**変形オイラー平均方程式**が得られる.

$$\frac{\partial \bar{u}}{\partial t} + \bar{v}^* \left[\frac{1}{a \cos \phi} \frac{\partial}{\partial \phi} (\bar{u} \cos \phi) - f \right] + \bar{w}^* \frac{\partial \bar{u}}{\partial z^*} - \bar{X} = \frac{1}{\rho_0 a \cos \phi} \nabla \cdot \mathbf{F}, \quad (\text{A.18a})$$

$$\bar{u} \left(f + \bar{u} \frac{\tan \phi}{a} \right) + \frac{1}{a} \frac{\partial \bar{\Phi}}{\partial \phi} = G. \quad (\text{A.18b})$$

$$\frac{\partial \bar{\Phi}}{\partial z^*} - \frac{Re^{-\kappa z^*/H}}{H} \bar{\theta} = 0. \quad (\text{A.18c})$$

$$\frac{1}{a \cos \phi} \left[\frac{\partial}{\partial \phi} (\bar{v}^* \cos \phi) \right] + \frac{1}{\rho_0} \frac{\partial}{\partial z^*} (\rho_0 \bar{w}^*) = 0. \quad (\text{A.18d})$$

$$\frac{\partial \bar{\theta}}{\partial t} + \frac{\bar{v}^*}{a} \frac{\partial \bar{\theta}}{\partial \phi} + \bar{w}^* \frac{\partial \bar{\theta}}{\partial z^*} - \bar{Q} = -\frac{1}{\rho_0} \frac{\partial}{\partial z^*} \left[\rho_0 \left(\bar{v}' \theta' \frac{\partial \bar{\theta}}{\partial \phi} + \bar{w}' \theta' \right) \right]. \quad (\text{A.18e})$$

関連図書

- [1] D.G. Andrews, J.R. Holton, and C.B. Leovy. Middle atmosphere dynamics, International Geophysics Series. Academic Press, 1987
- [2] J.R. Holton. The Dynamic Meteorology of the Stratosphere and Mesosphere, American Meteorological Society, 1975